

# Solid State Physics - A Summary

## 1 Crystallography

*Structure*: array of atoms, specified by a basis, plus lattice;

*Basis*: Group of atoms, which is repeated;

*Lattice*: Pattern of repeats;

*Primitive unit cell*: Cell with 1 lattice site only, and smallest area;

*Bravais Lattice*: Fill all space, and environment of all lattice sites identical;

*Packing fraction*: Fraction of space filled by touching spheres;

*Coordination number*: Number of nearest neighbours.

There are **5 2D Bravais lattices**: *square, rectangular, triangular, rhomic & oblique*.

There are **14 3D Bravais lattices**: we only consider a few *sc, bcc, fcc*.

Examples:

sc: Polonium

bcc: Li, Na, Fe

fcc: Sr, Ag, Cu

### 1.1 Miller Indices

Denoted  $(h, k, l)$ . To find the indices of a plane:

Get planes intercepts with axes. Reciprocate (putting 'bars' over negative intercepts). Multiply up, to integers.

We often refer to some structure, from the basis of some crystal:

Diamond: fcc lattice, with basis  $(0, 0, 0), (\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$ . Visualise this: at each fcc lattice site, put a C atom there, and one at  $(\frac{1}{4}, \frac{1}{4}, \frac{1}{4})$  relative to that site.

### 1.2 X-ray Diffraction

Probe crystals with  $\lambda$  of the order inter-site spacing.

We could use *thermal neutrons*, but very expensive to produce, but do scatter only once (which is good: easy to interpret data).

We can also use *electrons*, but results are very hard to interpret, because multiple scattering: use for surface investigations.

#### 1.2.1 Braggs Law

An X-ray is incident, at angle  $\theta$  to *plane*, which are spaced by  $d$ . Constructive interference if:

$$2d \sin \theta = n\lambda \tag{1.1}$$

For a plane with Miller indices  $h, k, l$ ; they are spaced by:

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}} \quad (1.2)$$

If  $a$  is the lattice spacing.

Intensity is higher for lower Miller indices: more atoms on planes.  
We get Bragg reflections for  $\sin \theta < 1$ . That is, for:

$$\frac{n\lambda}{2a} \sqrt{h^2 + k^2 + l^2} < 1 \quad (1.3)$$

**Indexing of lines:**

$$\frac{\sin \theta_2}{\sin \theta_1} = \frac{\sqrt{h_2^2 + k_2^2 + l_2^2}}{\sqrt{h_1^2 + k_1^2 + l_1^2}} \quad (1.4)$$

For reflections at  $\theta_1, \theta_2$ .

### 1.3 Reciprocal Lattice

If  $\mathbf{a}_i$  are the real space lattice basis vectors, and if  $\mathbf{b}_i$  are the reciprocal lattice basis vectors, then, we have the relations:

$$\mathbf{R} = n_i \mathbf{a}_i \quad \mathbf{G} = m_i \mathbf{b}_i \quad \mathbf{a}_i \cdot \mathbf{b}_j = 2\pi \delta_{ij} \quad e^{i\mathbf{R} \cdot \mathbf{G}} = 1 \quad (1.5)$$

We find  $\mathbf{b}_i$  via:

$$V \equiv \mathbf{a}_1 \cdot (\mathbf{a}_2 \times \mathbf{a}_3) \quad \mathbf{b}_1 = \frac{2\pi}{V} (\mathbf{a}_2 \times \mathbf{a}_3) \quad (1.6)$$

Cycle through indices to get other  $\mathbf{b}_i$ 's.

The spacing of planes, with Miller indices  $h, k, l$ ; is given by:

$$d = \frac{2\pi}{|\mathbf{G}|} \quad \mathbf{G} = h\mathbf{b}_1 + k\mathbf{b}_2 + l\mathbf{b}_3 \quad (1.7)$$

If a wave  $\mathbf{k}$  is incident to a plane, and  $\mathbf{k}'$  comes out; if *elastic scattering* then:

$$\mathbf{k} - \mathbf{k}' = \mathbf{G} \quad (1.8)$$

Thus, strong reflection for waves which differ by a reciprocal lattice vector (way to derive Bragg's law).

#### 1.3.1 Ewald Construction

Consider a circle, over lattice sites. Then strong reflection only if circle cuts a lattice site. Thus, low probability.

*Laue Method:* Vary the incident wavelength. Then, all lattice sites within the circle get picked up. We basically fire white X-rays at a single crystal, and observe the diffraction pattern. We see spots, and they have the same symmetry as the lattice.

*Powder Method:* Fire mono-chromatic X-rays through a hole in a cylinder of film. In the centre is a rod of specimen. Lines are seen in the film.

### 1.3.2 Effect of a Basis

The *structure factor* gives the amplitude of waves, from reflection:

$$S(\mathbf{G}) = \sum_{j=1}^N f_j(\mathbf{G}) e^{i\mathbf{G} \cdot \mathbf{r}_j} \quad (1.9)$$

For  $N$  basis vectors. Ignore the *atomic form factor*,  $f$  for mono-atomic lattices. It is related to the density of electrons.

Intensity of reflected wave: multiply by  $|S|^2$