

Laplace Transforms:

$$L\{f(t)\} \equiv \bar{f}(p) \equiv \int_0^{\infty} f(t)e^{-pt} dt$$

The operator is linear. That is:

$$\begin{aligned} L\{f_1(t) + f_2(t)\} &= L\{f_1(t)\} + L\{f_2(t)\} \\ L\{af(t)\} &= aL\{f(t)\} \end{aligned}$$

Assume the contribution from the infinite-limit is zero.

Some common transforms:

$$\begin{aligned} f(t) = 1 & \qquad \bar{f}(p) = \frac{1}{p} \\ f(t) = e^{at} & \qquad \bar{f}(p) = \frac{1}{p-a} \\ f(t) = t^n & \qquad \bar{f}(p) = \frac{n!}{p^{n+1}} \\ f(t) = t^n e^{at} & \qquad \bar{f}(p) = \frac{n!}{(p-a)^{n+1}} \end{aligned}$$

$$\begin{aligned} L\left\{\frac{dy}{dt}\right\} &= \int_0^{\infty} \frac{dy}{dt} e^{-pt} dt \\ &= \left[ye^{-pt}\right]_0^{\infty} + p \int_0^{\infty} ye^{-pt} dt \\ &= -y(0) + p\bar{y}(p) \\ L\left\{\frac{d^2y}{dt^2}\right\} &= \int_0^{\infty} \frac{d^2y}{dt^2} e^{-pt} dt \\ &= \left[\frac{dy}{dt} e^{-pt}\right]_0^{\infty} - \int_0^{\infty} \left(-p \frac{dy}{dt} e^{-pt}\right) dt \\ &= -y'(0) - py(0) + p^2 \bar{y}(p) \end{aligned}$$

Thus:

$$\begin{aligned} L\left\{\frac{dy}{dt}\right\} &= p\bar{y}(p) - y(0) \\ L\left\{\frac{d^2y}{dt^2}\right\} &= p^2 \bar{y}(p) - y'(0) - py(0) \end{aligned}$$

$$\begin{aligned} L\{e^{i\omega t}\} &= \int_0^{\infty} e^{i\omega t} e^{-pt} dt \\ &= \int_0^{\infty} e^{t(i\omega - p)} dt \\ &= -\frac{1}{i\omega - p} = \frac{1}{p - i\omega} \cdot \frac{p + i\omega}{p + i\omega} \\ &= \frac{p + i\omega}{p^2 + \omega^2} \end{aligned}$$

Thus, from this result, can see that:

$$\begin{aligned} L\{\cos \omega t\} &= L\{\operatorname{Re}\{e^{i\omega t}\}\} = \operatorname{Re}\{L\{e^{i\omega t}\}\} \\ &= \frac{p}{p^2 + \omega^2} \end{aligned}$$

And  $L\{\sin \omega t\} = \frac{\omega}{p^2 + \omega^2}$

To invert, if we cannot do by inspection, use the inversion integral:

$$f(t) = L^{-1}\{\bar{f}(p)\} = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} \bar{f}(p) e^{pt} dp$$

Where  $c$  is chosen to the right of all singularities of  $\bar{f}(p)e^{pt}$ .

c.f. Jordan's Lemma.

Thus, only get contributions to integral from singularities & their residues:

$$f(t) = \frac{1}{2\pi i} \left[ 2\pi i \sum_{\text{to left of } c} \operatorname{Res}\{\bar{f}(p)e^{pt}\} \right]$$

We can use Laplace transforms to solve differential equations:

$$\frac{dy}{dt} + y = e^{-2t} \quad y(0) = 1$$

Take LT of whole thing, remembering that it is linear:

$$L\left\{\frac{dy}{dt}\right\} + L\{y(t)\} = L\{e^{-2t}\}$$

Or:

$$p\bar{y}(p) - \underbrace{y(0)}_{=0} + \bar{y}(p) = \frac{1}{p+2} \quad \text{making use of initial conditions}$$

$$\bar{y}(p) = \frac{p+3}{(p+2)(p+1)} = \frac{2}{p+1} - \frac{1}{p+2}$$

Which can be inverted by inspection:

$$y(t) = 2e^{-t} - e^{-2t}$$

<u>Function</u>	<u>Laplace Transform</u>
$e^{at}$	$\frac{1}{p-a}$
$t^n$	$\frac{n!}{p^{n+1}}$
$t^n e^{at}$	$\frac{n!}{(p-a)^{n+1}}$
$\cos \omega t$	$\frac{p}{p^2 + \omega^2}$
$\sin \omega t$	$\frac{\omega}{p^2 + \omega^2}$
$\frac{dy}{dt}$	$p\bar{y}(p) - y(0)$
$\frac{d^2 y}{dt^2}$	$p^2 \bar{y}(p) - y'(0) - py(0)$